

Solution of an Initial Value Problem of Cauchy Type for One Equation

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Abstract: In this work, using pre-constructed partial solutions, a representation of the solution to the Cauchy problem for an equation with two variables, in which both derivatives are of fractional order, is obtained.

Keywords: RiemannLiouville fractional derivative, self-similar solution, Cauchy problem, explicit solution.

1 Introduction and Finding Particular Solutions

Consider the equation

$$L[u] \equiv D_{0y}^\alpha u(x, y) - dD_{0x}^\beta u(x, y) = 0, \tag{1}$$

where

$$x, y > 0, 0 < \alpha < 1; 1 < \beta < 2; \alpha + \beta < 2; d < 0,$$

and $D_{0y}^\alpha, D_{0x}^\beta$ are derivatives of fractional order in the sense of Riemann-Liouville, respectively, orders α, β (see [1]) :

$$D_{st}^\nu \varphi(t) = \begin{cases} \frac{\text{sign}(t-s)}{\Gamma(-\nu)} \int_s^t \frac{\varphi(\tau) d\tau}{|t-\tau|^{\nu+1}}, \nu < 0, \\ \varphi(t), \nu = 0, \\ \text{sign}^n(t-s) \frac{d^n}{dt^n} D_{st}^{\nu-n} \varphi(t), n-1 < \nu \leq n. \end{cases}$$

Recently, specialists have increasingly intensively studied equations that have a fractional order. This is due to the wide application of these equations in the natural sciences and in life (see, for example, [1]-[6] and others).

When solving initial and boundary problems, it is important to know the fundamental solution. One of the methods for finding a fundamental solution is the method of preliminary construction of a self-similar solution. For example, this method was used to construct a fundamental solution of the heat equation [7]. The construction of self-similar solutions themselves is also important from both theoretical and practical points of view. In the articles [8]-[9] particular solutions of the type of self-similar solutions for model equations of high integer order were found. In works [10]-[15] self-similar solutions were constructed for equations with fractional derivatives using special integro-differential operators.

In this part of the article we will construct self-similar solutions to equation (1) in the form of the following series:

$$u(x, y) = y^b \sum_{n=0}^{\infty} c_n (x^a y^c)^{n+\gamma} = \sum_{n=0}^{\infty} c_n x^{an+a\gamma} y^{cn+c\gamma+b}, \tag{2}$$

where the parameters a, b, c, γ are unknown yet. Let be

$$\overline{(a)}_s = a(a-1) \dots (a-(s-1)), \overline{(a)}_0 = 1, \overline{(a)}_1 = a.$$

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Taking (2) into account, we formally have that

$$\begin{aligned} D_{0y}^{\alpha} u(x, y) &= \frac{1}{\Gamma(1-\alpha)} \sum_{n=0}^{\infty} c_n x^{an+a\gamma} \frac{d}{dy} \int_0^y \frac{\tau^{cn+c\gamma+b} d\tau}{(y-\tau)^{\alpha}} = \\ &= \sum_{n=0}^{\infty} c_n (x^a y^c)^{n+\gamma} y^{b-\alpha} \frac{\Gamma(cn+c\gamma+b+1)}{\Gamma(cn+c\gamma+b-\alpha+1)}, \end{aligned} \quad (3)$$

the same way

$$D_{0x}^{\beta} u(x, y) = \sum_{n=0}^{\infty} \frac{1}{(an+a\gamma+2-\beta)_2} c_n (x^a y^c)^{n+\gamma} x^{-\beta} y^b \frac{\Gamma(an+a\gamma+1)}{\Gamma(an+a\gamma+3-\beta)}. \quad (4)$$

Let now $a = \beta$, $c = -\alpha$, then taking into account (3),(4) from equation (1) we get

$$\begin{aligned} &\sum_{n=0}^{\infty} c_n (x^a y^c)^n x^{\beta} y^{-\alpha} \frac{\Gamma(cn+c\gamma+b+1)}{\Gamma(cn+c\gamma+b-\alpha+1)} = \\ &= d \sum_{n=0}^{\infty} \frac{1}{(an+a\gamma+2-\beta)_2} c_n (x^a y^c)^n \frac{\Gamma(an+a\gamma+1)}{\Gamma(an+a\gamma+3-\beta)}, \end{aligned}$$

or

$$\begin{aligned} &\sum_{n=0}^{\infty} c_n (x^{\beta} y^{-\alpha})^{n+1} \frac{\Gamma(-\alpha(n+\gamma)+b+1)}{\Gamma(-\alpha(n+\gamma+1)+b+1)} = \\ &= d \sum_{n=0}^{\infty} \frac{1}{(\beta n + \beta\gamma + 2 - \beta)_2} c_n (x^a y^c)^n \frac{\Gamma(an+a\gamma+1)}{\Gamma(\beta n + \beta\gamma + 3 - \beta)}, \end{aligned}$$

hence the equality must be satisfied that

$$\gamma_1 = 1 - \frac{1}{\beta}, \gamma_2 = 1 - \frac{2}{\beta}.$$

Taking this into account, after some calculations and transformations we obtain a formula for finding the coefficients c_n in the form:

$$\begin{aligned} c_0 &= \frac{1}{\Gamma(-\alpha\gamma+b+1)\Gamma(\beta\gamma+1)}, j = 1, 2; \\ c_n &= \frac{1}{d^n \Gamma(-\alpha n - \alpha + \frac{\alpha}{\beta} j + b + 1) \Gamma(\beta n + \beta - j + 1)}, j = 1, 2; n = 1, 2, \dots \end{aligned}$$

So, we have obtained the following formal partial solutions of equation (1), such as self-similar solutions, in the following form:

$$u_j(x, y) = x^{\beta-j} y^{-\alpha+j\frac{\alpha}{\beta}+b} e_{\beta,\alpha}^{\beta-j+1, -\alpha+j\frac{\alpha}{\beta}+b+1} \left(\frac{1}{d} x^{\beta} y^{-\alpha} \right), j = 1, 2, \quad (5)$$

where

$$e_{\lambda,\nu}^{\mu,\delta}(z) = \sum_{n=0}^{\infty} \frac{z^n}{\Gamma(\lambda n + \mu) \Gamma(\delta - \nu n)}, \lambda > \nu, \lambda > 0, z \in \mathbb{C}$$

is a Wright type function [1]. All the above calculations were of a formal nature. Let us give expression (5) a legitimate meaning. The following theorem is true.

Theorem 1. If $d < 0$; $x > 0, y > 0$; $0 < \alpha < 1$; $1 < \beta < 2$; $\alpha + \beta < 2$; $j\frac{\alpha}{\beta} + b + 1 > 0$, $j = 1, 2$; then the expressions (5) are partial solutions of the equation (1).

Proof. By direct calculation, we have

$$D_{0x}^{\beta} u_j(x, y) = \frac{1}{d} x^{\beta-j} y^{b-2\alpha+j\frac{\alpha}{\beta}} e_{\beta,\alpha}^{-j+1+\beta, -2\alpha+\frac{\alpha}{\beta}j+b+1} \left(\frac{1}{d} x^{\beta} y^{-\alpha} \right).$$

Using formula (2.2.12) from [1], we obtain

$$x^{\beta-j} D_{0y}^{\beta} \left(y^{-\alpha+j\frac{\alpha}{\beta}+b} e_{\beta,\alpha}^{\beta-j+1, -\alpha+j\frac{\alpha}{\beta}+b+1} \left(\frac{1}{d} x^{\beta} y^{-\alpha} \right) \right) =$$

$$= x^{\beta-j} y^{-2\alpha+j\frac{\alpha}{\beta}+b} e_{\beta,\alpha}^{\beta-j+1,-2\alpha+j\frac{\alpha}{\beta}+b+1} \left(\frac{1}{d} x^{\beta} y^{-\alpha} \right).$$

Substituting the found fractional derivatives into the equation (1), we obtain an identity.

Theorem 1 is proved.

Solutions of the form (5) coincide with the solutions obtained in [11]. Let now

$$b = \alpha - 1 - \frac{\alpha}{\beta}, d = -1,$$

consider expressions

$$V_1(x, \xi, y, \eta) = |x - \xi|^{\beta-1} (y - \eta)^{-1} e_{\beta,\alpha}^{\beta,0} \left(-|x - \xi|^{\beta} (y - \eta)^{-\alpha} \right),$$

$$V_2 = |x - \xi|^{\beta-2} (y - \eta)^{\frac{\alpha}{\beta}-1} e_{\beta,\alpha}^{\beta-1,\frac{\alpha}{\beta}} \left(-|x - \xi|^{\beta} (y - \eta)^{-\alpha} \right),$$

known estimate [1]

$$|V_1(x, \xi, y, \eta)| \leq C |x - \xi|^{\beta-1-\beta\theta} |y - \eta|^{-1+\alpha\theta}, C - const, 0 \leq \theta \leq 2,$$

$$|V_2(x, \xi, y, \eta)| \leq C |x - \xi|^{\beta-2-\beta\theta} (y - \eta)^{\frac{\alpha}{\beta}-1+\alpha\theta}.$$

True lemma.

Lemma 1. Let $\varphi(x) \in C(R)$ and $|\varphi(x)| \leq S, \forall x \in R, 0 < S - const$, then

$$\lim_{y \rightarrow +0} \left(y^{1-\alpha} \int_{-\infty}^{+\infty} V_1(x, \xi, y, 0) \varphi(\xi) d\xi \right) = 2 \frac{\varphi(x)}{\Gamma(\alpha)}.$$

Proof. We have

$$y^{1-\alpha} \int_{-\infty}^{+\infty} V_1(x, \xi, y, 0) \varphi(\xi) d\xi = I_1(x, y) + I_2(x, y),$$

where

$$I_1 = y^{1-\alpha} \int_{-\infty}^x V_1(x, \xi, y, 0) \varphi(\xi) d\xi,$$

$$I_2 = y^{1-\alpha} \int_x^{+\infty} V_1(x, \xi, y, 0) \varphi(\xi) d\xi.$$

Let's make a change of variables

$$z = (x - \xi)^{\beta} y^{-\alpha}, \xi = x - z^{\frac{1}{\beta}} y^{\frac{\alpha}{\beta}},$$

then we will have

$$I_1 = \frac{\varphi(x)}{\beta} \int_0^{+\infty} e_{\beta,\alpha}^{\beta,0}(-z) dz + \frac{1}{\beta} \int_0^{+\infty} e_{\beta,\alpha}^{\beta,0}(-z) \left(\varphi \left(x - z^{\frac{1}{\beta}} y^{\frac{\alpha}{\beta}} \right) - \varphi(x) \right) dz.$$

Let's consider each term separately. Considering formula (2.2.7) from [1], we have

$$\frac{\varphi(x)}{\beta} \int_0^{+\infty} e_{\beta,\alpha}^{\beta,0}(-z) dz = \frac{\varphi(x)}{\beta} \lim_{A \rightarrow +\infty} \sum_{n=0}^{\infty} \int_0^A \frac{(-1)^n z^n dz}{\Gamma(\beta n + \beta) \Gamma(-\alpha n)} =$$

$$= \varphi(x) \lim_{A \rightarrow +\infty} \sum_{n=0}^{\infty} \frac{(-1)^n A^{n+1}}{(\beta n + \beta) \Gamma(\beta n + \beta) \Gamma(-\alpha n)} =$$

$$= \varphi(x) \lim_{A \rightarrow +\infty} A e_{\beta,\alpha}^{\beta+1,0}(-A) = \frac{\varphi(x)}{\Gamma(\alpha)}. \tag{6}$$

Further

$$\begin{aligned}
 \int_0^{+\infty} e_{\beta,\alpha}^{\beta,0}(-z) \left(\varphi \left(x - z^{\frac{1}{\beta}} y^{\frac{\alpha}{\beta}} \right) - \varphi(x) \right) dz &= \int_0^M e_{\beta,\alpha}^{\beta,0}(-z) \left(\varphi \left(x - z^{\frac{1}{\beta}} y^{\frac{\alpha}{\beta}} \right) - \varphi(x) \right) dz + \\
 &+ \int_M^{+\infty} e_{\beta,\alpha}^{\beta,0}(-z) \left(\varphi \left(x - z^{\frac{1}{\beta}} y^{\frac{\alpha}{\beta}} \right) - \varphi(x) \right) dz, \\
 \forall \varepsilon > 0, \exists M > 0, 0 < y < \delta &\Rightarrow \\
 \int_0^M \left| e_{\beta,\alpha}^{\beta,0}(-z) \right| \left| \varphi \left(x - z^{\frac{1}{\beta}} y^{\frac{\alpha}{\beta}} \right) - \varphi(x) \right| dz &\leq \\
 \leq \sup_{0 < z < M, 0 < y < \delta} \left| \varphi \left(x - z^{\frac{1}{\beta}} y^{\frac{\alpha}{\beta}} \right) - \varphi(x) \right| \int_0^M \left| e_{\beta,\alpha}^{\beta,0}(-z) \right| dz &\leq \frac{\varepsilon}{2}, \\
 \int_M^{+\infty} \left| e_{\beta,\alpha}^{\beta,0}(-z) \right| \left| \varphi \left(x - z^{\frac{1}{\beta}} y^{\frac{\alpha}{\beta}} \right) - \varphi(x) \right| dz &\leq \\
 \leq 2S \int_M^{+\infty} \left| e_{\beta,\alpha}^{\beta,0}(-z) \right| dz = 2S \int_M^{+\infty} z^{-2} \left| z^2 e_{\beta,\alpha}^{\beta,0}(-z) \right| dz,
 \end{aligned} \tag{7}$$

taking into account relation (2.2.8) from [1], we have

$$\lim_{z \rightarrow +\infty} \left(z^2 e_{\beta,\alpha}^{\beta,0}(-z) \right) = \frac{1}{\Gamma(-\beta)\Gamma(2\alpha)},$$

from here

$$\int_M^{+\infty} \left| e_{\beta,\alpha}^{\beta,0}(-z) \right| \left| \varphi \left(x - z^{\frac{1}{\beta}} y^{\frac{\alpha}{\beta}} \right) - \varphi(x) \right| dz \leq \frac{\varepsilon}{2}. \tag{8}$$

From (6),(7) and (8) we get

$$\lim_{y \rightarrow +0} I_1(x,y) = \frac{\varphi(x)}{\Gamma(\alpha)}.$$

Similarly

$$\lim_{y \rightarrow +0} I_2(x,y) = \frac{\varphi(x)}{\Gamma(\alpha)}.$$

Lemma 1 is proved.

Lemma 2. Let be $\varphi(x) \in C(R)$ and $|\varphi(x)| \leq M, \forall x \in R, 0 < M - const$, then, function

$$u(x,y) = \int_{-\infty}^{+\infty} V_1(x,\xi,y,0) \varphi(\xi) d\xi, \tag{9}$$

is a solution to equation (1).

Proof. According to the estimate of function V_1 , the improper integral (9) converges. Further, we have

$$\begin{aligned}
 D_{0x}^\alpha \left\{ \int_{-\infty}^{+\infty} V_1(x,\xi,y,0) \varphi(\xi) d\xi \right\} &= \\
 = \frac{1}{y} \frac{1}{\Gamma(2-\alpha)} \frac{\partial^2}{\partial x^2} \int_0^x \left(\int_{-\infty}^{+\infty} \frac{|t-\xi|^{\beta-1} e_{\beta,\alpha}^{\beta,0}(-|t-\xi|^\beta y^{-\alpha}) \varphi(\xi)}{|x-t|^{\alpha-1}} d\xi \right) dt &=
 \end{aligned}$$

$$\begin{aligned}
 &= \frac{1}{y} \frac{1}{\Gamma(2-\alpha)} \frac{\partial^2}{\partial x^2} \int_{-\infty}^{+\infty} \varphi(\xi) d\xi \left(\int_0^x \frac{|t-\xi|^{\beta-1} e_{\beta,\alpha}^{\beta,0}(-|t-\xi|^\beta y^{-\alpha})}{|x-t|^{\alpha-1}} dt \right) = \\
 &= \int_{-\infty}^{+\infty} \varphi(\xi) D_{0x}^\alpha V_1(x, \xi, y, 0) d\xi.
 \end{aligned}$$

The convergence of the last improper integral is proved in the same way as in Lemma 1. Similarly,

$$D_{0y}^\beta \left\{ \int_{-\infty}^{+\infty} V_1(x, \xi, y, 0) \varphi(\xi) d\xi \right\} = \int_{-\infty}^{+\infty} \varphi(\xi) D_{0y}^\beta V_1(x, \xi, y, 0) d\xi.$$

From here

$$D_{0y}^\alpha u(x, y) + D_{0x}^\beta u(x, y) = 0.$$

Lemma 2 is proved.

2 Representation of the Solution to the Initial Problem

In this part of the article, we will apply the constructed self-similar solutions to solve the following initial value problem, such as the Cauchy problem.

Cauchy problem. In the region $\Omega = \{(x, y) : -\infty < x < +\infty, 0 < y \leq T\}, 0 < T < +\infty, T - const$, find a solution to the problem :

$$\begin{cases} D_{0y}^\alpha u(x, y) + D_{0x}^\beta u(x, y) = 0, \\ \lim_{y \rightarrow +0} (y^{\alpha-1} u(x, y)) = \varphi(x), \end{cases}$$

where

$$\begin{aligned}
 &y^{\alpha-1} u(x, y) \in C(\overline{\Omega}), D_{0y}^\alpha u(x, y), D_{0x}^\beta u(x, y) \in C(\Omega), \\
 &0 < \alpha < 1; 1 < \beta < 2; \alpha + \beta < 2, \\
 &\varphi(x) \in C(R), |\varphi(x)| < M < \infty, \forall x \in R, 0 < M - const.
 \end{aligned}$$

The results of Lemma 1 and Lemma 2 imply the validity of the following theorem.

Theorem 2. Function having the form:

$$u(x, y) = \frac{\Gamma(\alpha)}{2} \int_{-\infty}^{+\infty} \frac{|x-\xi|^{\beta-1}}{y} e_{\beta,\alpha}^{\beta,0} \left(-\frac{|x-\xi|^\beta}{y^\alpha} \right) \varphi(\xi) d\xi$$

is a solution to the Cauchy problem.

Note. Note that the constructed self-similar solutions can be used to obtain a fundamental solution of the heat equation. Indeed, the heat equation

$$u_y(x, y) - u_{xx}(x, y) = 0,$$

can be written in equation form (1) (limiting case $\alpha = 1, \beta = 2$) as

$$u_y(x, y) + u_{xx}(ix, y) = 0,$$

further, having made some transformations, we will have

$$\begin{aligned}
 V_2 &= \frac{1}{\sqrt{y-\eta}} e_{2,1}^{1,\frac{1}{2}} \left(-i^2 |x-\xi|^2 (y-\eta)^{-1} \right) = \\
 &= \frac{1}{\sqrt{y-\eta}} \sum_{n=0}^{\infty} \frac{\left(\frac{(x-\xi)^2}{y-\eta} \right)^n}{\Gamma(2n+1) \Gamma(-n+\frac{1}{2})} = \frac{1}{\sqrt{\pi(y-\eta)}} e^{-\frac{(x-\xi)^2}{4(y-\eta)}}.
 \end{aligned}$$

That is, have obtained a fundamental solution of the heat equation.

References

- [1] A. V. Pskhu , *Uravneniya v chastnykh proizvodnykh drobnogo poryadka [Partial differential equations of fractional order]*, Moscow: NAUKA, 2005. (in Russian).
 - [2] J. Singh, Analysis of fractional blood alcohol model with composite fractional derivative, *Chaos Soliton. Fract.* **140**, 110127 (2020).
 - [3] J. Singh , D. Kumar and D. Baleanu, A new analysis of fractional fish farm model associated with Mittag-Leffler-type kernel, *Int. J. Biomath.* **13** (2), 2050010 (2020).
 - [4] V. V. Uchaikin, *Fractional derivatives for physicists and engineers*, Volume I Background and Theory Volume II Applications, Springer, 2013.
 - [5] R. R. Niqmatullin, Theoretical explanation of the "universal response", *Phys. Stat. Sol. (b)* **123** (2), 739–745 (1984).
 - [6] R. R. Niqmatullin, The realization of the generalized transfer equation in a medium with fractal geometry, *Phys. Stat. Sol. (b)* **133** (1), 425–430 (1986).
 - [7] A. N. Tikhonov and A. A. Samarskiy , *The equations of mathematical physics*, 2nd ed., Gosudarstv. Izdat. Tehn.-Teor. Lit., Moscow, 1953.(in Russian).
 - [8] B. Yu. Irgashev, Construction of singular particular solutions expressed via hypergeometric functions for an equation with multiple characteristics, *Differ. Equ.* **56** (3), 315–323 (2020).
 - [9] M. Ruzhansky and A. Hasanov, Self similar solutions of some model degenerate partial differential equations of the second, third, and fourth order, *Lobachevskii J. Math.* **41** (6), 1103–1114 (2020).
 - [10] E. Buckwar and Y. Luchko, Invariance of a partial differential equation of fractional order under the Lie group of scaling transformations, *J. Math. Anal. App.* **227**(1), 81–97 (1998).
 - [11] Y. Luchko and R. Gorenflo, Scaleinvariant solutions of a partial differential equation of fractional order, *Fract. Calc. Appl. Anal.* **1** (1), 63–78 (1998).
 - [12] J. S. Duan , A. P. Guo and W. Z. Yun, Similarity solution for fractional diffusion equation, *Abst. Appl. Anal.* Article ID 548126, 1–5 (2014).
 - [13] B. Basti and N. Benhamidouche, Existence results of self-similar solutions to the Caputo-types space-fractional heat equation, *Survey. Math. App.* **15**, 153–168 (2020).
 - [14] R. Gorenflo , Y. Luchko and F. Mainardi, Wright functions as scale-invariant solutions of the diffusion-wave equation, *J. Comput. Appl. Math.* **118** (1-2) 175–191 (2000).
 - [15] F. Al-Musalhi and E. Karimov, On self-similar solutions of time and space fractional sub-diffusion equations, *An Int. J. Opt. Cont. Theor. Appl.* **11**(3), 16–27 (2021).
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