

Fekete-Szegő Inequality and Application of Poisson Distribution Series for some Subclasses of Analytic Functions related with Bessel Functions

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Abstract: The determination of this current paper is to find certain coefficient estimates, Fekete-Szegő inequality results for a normalized analytic function defined in the open unit disk $\mathbb{D} = \{z : z \in \mathbb{C} \text{ and } |z| < 1\}$ by convolution operator with Bessel function. In particular, we derived Fekete-Szegő inequality for a class of functions defined through Poisson distribution. The results presented in this paper would generalize some related works of several earlier authors.

Keywords: Analytic functions, Starlike functions, Convex functions, Subordination, Fekete-Szegő inequality, Poisson distribution, Bessel function, Hadamard product.

1 Introduction, Definitions and Preliminary results

Let \mathcal{A} be the set of all analytic functions, comprising of the functions

$$f(z) = z + \sum_{n=2}^{\infty} a_n z^n, \tag{1}$$

in the open unit disc

$$\mathbb{D} := \{z : z \in \mathbb{C} \text{ and } |z| < 1\}$$

and \mathcal{S} be the subclass of \mathcal{A} comprising of univalent functions. Let the functions f and g be analytic in \mathbb{D} . We say that the function f is subordinate to g , if there exists a Schwarz function ω , which is analytic in \mathbb{D} with

$$\omega(0) = 0 \quad \text{and} \quad |\omega(z)| < 1 \quad (z \in \mathbb{D}),$$

such that

$$f(z) = g(\omega(z)).$$

This subordination is denoted by

$$f \prec g \quad \text{or} \quad f(z) \prec g(z) \quad (z \in \mathbb{D}).$$

It is well known that (see [1]), if the function g is univalent in \mathbb{D} , then

$$f \prec g \quad (z \in \mathbb{D}) \iff f(0) = g(0) \quad \text{and} \quad f(\mathbb{D}) \subseteq g(\mathbb{D}).$$

We recall here a Bessel function of the first kind of order ν , denoted by $J_\nu(z)$, is defined by the infinite series:

$$J_\nu(z) = \sum_{n=0}^{\infty} \frac{(-1)^n}{n! \Gamma(n + \nu + 1)} \left(\frac{z}{2}\right)^{2n+\nu} \quad (z \in \mathbb{C}, \nu \in \mathbb{R}) \tag{2}$$

which is the particular solution of the second order linear homogeneous differential equation

$$z^2 \omega''(z) + z\omega'(z) + [z^2 - \nu^2]\omega(z) = 0, \tag{3}$$

where $\nu \in \mathbb{C}$, which is the natural Bessel's equation. Solutions of (3) are referred to as Bessel function of order ν . Although the series defined in (2) is convergent every where, in general J_ν is not univalent in \mathbb{U} . Latterly, Szász and Kupán[2] inspected the univalence of the normalized Bessel function of the first kind

$$u_\nu : \mathbb{D} \rightarrow \mathbb{C}$$

given by the transformation (see also[3,4])

$$u_\nu(z) = 2^\nu \Gamma(\nu + 1) z^{1-\frac{\nu}{2}} J_\nu(\sqrt{z}), \quad \sqrt{1} = 1.$$

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We can express $u_v(z)$ as

$$u_v(z) = z + \sum_{n=1}^{\infty} \frac{(-1/4)^{n-1} \Gamma(v+1)}{(n-1)! \Gamma(n+v)} z^n, \quad (4)$$

is analytic on \mathbb{C} and satisfies the differential equation

$$4z^2 u''(z) + 4(v+1)zu'(z) + cu(z) = 0.$$

Quantum calculus (q -calculus and h -calculus) simply the study of classical calculus without the notion of limits. Here, h represents Planck's constant, while q represents quantum. Due to its application in a variety of branches such as physics, mathematics, the area of q -calculus has added excessive prominence for researchers. The first study on q -calculus was systematically established by Jackson [5]. Recently, Kanas and Raducanu [6] defined a q -analogue of the Ruscheweyh differential operator [7] by using the concept of convolution and then studied some of its properties by Aldweby and Darus [8]. Now, we give some notational details of q -calculus which are used in the paper.

For $f \in \mathcal{A}$ the Jackson's q -derivative ($0 < q < 1$) is expressed by

$$\mathfrak{D}_q f(z) = \begin{cases} \frac{f(z) - f(qz)}{(1-q)z}, & z \neq 0 \\ f'(0), & z = 0 \end{cases} \quad (5)$$

and $\mathfrak{D}_q^2 f(z) = \mathfrak{D}_q(\mathfrak{D}_q f(z))$. Thus, from (5), we presume that

$$\mathfrak{D}_q f(z) = 1 + \sum_{n=2}^{\infty} [n]_q a_n z^{n-1},$$

where

$$[n]_q = \frac{1 - q^n}{1 - q}.$$

If $q \rightarrow 1^-$, we get $[n]_q \rightarrow n$. For the function $h(z) = z^n$, we get $\mathfrak{D}_q h(z) = \mathfrak{D}_q z^n = \frac{1 - q^n}{1 - q} z^{n-1} = [n]_q z^{n-1}$ and $\lim_{q \rightarrow 1^-} \mathfrak{D}_q h(z) = \lim_{q \rightarrow 1^-} ([n]_q z^{n-1}) = n z^{n-1} = h'(z)$, where h' is the usual derivative.

For $0 < q < 1$, and $f \in \mathcal{A}$ of the form (1), the q -derivative of u_v is defined by:

$$\begin{aligned} \mathfrak{D}_q u_v(z) &= \mathfrak{D}_q \left[z + \sum_{n=1}^{\infty} \frac{(-1/4)^{n-1} \Gamma(v+1)}{(n-1)! \Gamma(n+v)} z^n \right] \\ &= \frac{u_v(z) - u_v(qz)}{(1-q)z} \\ &= 1 + \sum_{n=1}^{\infty} \frac{(-1/4)^{n-1} \Gamma(v+1)}{(n-1)! \Gamma(n+v)} [n]_q z^{n-1}. \end{aligned} \quad (6)$$

Let $a \in \mathbb{R}$ and $n \in \mathbb{N}$. The q -generalized Pochhammer symbol is defined by

$$[a; n]_q = [a]_q [a+1]_q [a+2]_q \dots [a+n-1]_q \quad (7)$$

and for $a > 0$ the q -gamma function is defined by

$$\Gamma_q(a+1) = [a]_q \Gamma_q(a) \quad \text{and} \quad \Gamma_q(1) = 1. \quad (8)$$

For $f \in \mathcal{A}$, Kanas and Raducanu [6] defined the Ruscheweyh q -differential operator as below:

$$\mathcal{R}_q^\delta f(z) = f(z) * F_{q, \delta+1}(z) \quad (\delta > -1, z \in \mathbb{U}) \quad (9)$$

where

$$\begin{aligned} F_{q, \delta+1}(z) &= z + \sum_{n=2}^{\infty} \frac{\Gamma_q(n+\delta)}{[n-1]_q! \Gamma_q(1+\delta)} z^n \\ &= z + \sum_{n=2}^{\infty} \frac{[\delta+1; n]_q}{[n-1]_q!} z^n. \end{aligned} \quad (10)$$

using, (9) and (10), Aldweby and Darus [8] defined the q -analogue of Ruscheweyh operator $\mathcal{R}_q^\delta : \mathcal{A} \rightarrow \mathcal{A}$ as follows:

$$\mathcal{R}_q^\delta f(z) = z + \sum_{n=2}^{\infty} \frac{\Gamma_q(n+\delta)}{[n-1]_q! \Gamma_q(1+\delta)} a_n z^n \quad (z \in \mathbb{U}).. \quad (11)$$

As $q \rightarrow 1^-$, we note that

$$\mathcal{R}_q^0 f(z) = f(z), \quad \mathcal{R}_q^1 f(z) = z f'(z),$$

It is easy to check that

$$z \mathfrak{D}_q (F_{q, \delta+1}(z)) = \left(1 + \frac{[\delta]}{q^\delta} \right) F_{q, \delta+2}(z) - \frac{[\delta]}{q^\delta} F_{q, \delta+1}(z), \quad (12)$$

$z \in \mathbb{U}$. From (9), (12) and by the concept of Hadamard product, we have

$$z (\mathcal{R}_q^\delta f(z))' = (1 + \delta) \mathcal{R}_q^{1+\delta} f(z) - \delta \mathcal{R}_q^\delta f(z), \quad (z \in \mathbb{U}). \quad (13)$$

From (11), as $q \rightarrow 1^-$ we note that

$$\begin{aligned} \lim_{q \rightarrow 1^-} F_{q, \delta+1}(z) &= \frac{z}{(1-z)^{\delta+1}}, \\ \lim_{q \rightarrow 1^-} \mathcal{R}_q^\delta f(z) &= f(z) * \frac{z}{(1-z)^{\delta+1}} \end{aligned}$$

the usual Ruscheweyh derivative [7].

By the description of q -derivative and the perception of Hadamard product, we describe the linear operator

$$\Psi_{v,q}^\delta : \mathcal{A} \rightarrow \mathcal{A}$$

defined by

$$\begin{aligned} \Psi_{v,q}^\delta f(z) &= u_v(z) * (\mathcal{R}_q^\delta f(z)) \\ &= z + \sum_{n=2}^{\infty} \frac{(-1/4)^{n-1} \Gamma(v+1)}{(n-1)! \Gamma(n+v)} \frac{\Gamma_q(n+\delta)}{[n-1]_q! \Gamma_q(1+\delta)} a_n z^n \end{aligned} \quad (14)$$

$$= z + \sum_{n=2}^{\infty} \Theta_n a_n z^n \quad (z \in \mathbb{U}) \quad (16)$$

where

$$\Theta_n = \frac{(-1/4)^{n-1} \Gamma(v+1)}{(n-1)! \Gamma(n+v)} \frac{\Gamma_q(n+\delta)}{[n-1]_q! \Gamma_q(1+\delta)}. \quad (17)$$

Ma and Minda[9], unified various subclasses of starlike and convex functions for which either of the functions

$$\frac{zf'(z)}{f(z)} \quad \text{and} \quad 1 + \frac{zf''(z)}{f'(z)} \quad (z \in \mathbb{D}).$$

is subordinate to a more general superordinate function and denoted such function classes by $\mathcal{S}^*(\phi)$ and $\mathcal{C}(\Phi)$, respectively. For this purpose, they considered an analytic function Φ as below ,

Definition 1.[9] Suppose Φ is an analytic function such that

1. $\Re(\Phi) > 0$ in \mathbb{U}
2. $\Phi(0) = 1, \quad \Phi'(0) > 0$
3. Φ maps \mathbb{U} onto a region starlike with respect to 1 and symmetric with respect to the real axis.

Further they gave $\Phi(z)$ in series by

$$\Phi(z) = 1 + B_1z + B_2z^2 + B_3z^3 + \dots, \quad (18)$$

where B_n 's are real with $B_1 > 0; B_2 \geq 0$.

Fixing $\Phi(z) = \frac{1 + (1 - 2\beta)z}{1 - z}, \quad (z \in \mathbb{D}; 0 \leq \beta < 1)$

we get the well-known classes $\mathcal{S}^*(\beta)$ (and $\mathcal{C}(\beta)$) of starlike functions (and the class of convex functions) of order $\beta(0 \leq \beta < 1)$ respectively. In [10], Guo and Liu defined a subclass $M(\mu, \lambda, \rho)$ as below which unifies certain subclasses of analytic functions.

Let $\mu \geq 0, \lambda \geq 0$ and $0 \leq \rho < 1$ and $f \in \mathcal{A}$. We say that $f \in M(\mu, \lambda, \rho)$ if it hold the analytic criterion

$$\Re \left\{ \frac{zf'(z)}{f(z)} \left(\frac{f(z)}{z} \right)^\mu + \lambda \left[1 + \frac{zf''(z)}{f'(z)} - \frac{zf'(z)}{f(z)} + \mu \left(\frac{zf'(z)}{f(z)} - 1 \right) \right] \right\} > \rho.$$

In [9], the authors have obtained the Fekete-Szegő inequality f in $\mathcal{S}^*(\Phi)$ and $\mathcal{C}(\Phi)$. For a brief history of Fekete-Szegő problem various subclasses of analytic functions, one may refer to [11] and the references cited there in. Motivated essentially by the aforementioned works on Fekete-Szegő inequality and the definition of hadamard product we define a more general class of analytic functions which unifies the class $\mathcal{S}^*(\Phi)$ and $\mathcal{C}(\Phi)$, $\mathcal{M}_\lambda(\Phi)$ based on Bessel function. Also, we give applications of our results to certain functions defined through of Poisson distribution series.

Now, we define the following new function class $\mathcal{G}_{\mu, \lambda}^{v, \delta, q}(\Phi)$:

Definition 2. For $\mu \geq 0, \lambda \geq 0$, let $\Phi(z)$ be in Definition 1 and $f \in \mathcal{A}$ is in the class $\mathcal{G}_{\mu, \lambda}^{v, \delta, q}(\Phi)$ if

$$\frac{z(\Psi_{v,q}^\delta f(z))'}{\Psi_{v,q}^\delta f(z)} \left(\frac{\Psi_{v,q}^\delta f(z)}{z} \right)^\mu + \lambda \left[1 + \frac{z(\Psi_{v,q}^\delta f(z))''}{(\Psi_{v,q}^\delta f(z))'} - \frac{z(\Psi_{v,q}^\delta f(z))'}{\Psi_{v,q}^\delta f(z)} \right] + \mu \left(\frac{z(\Psi_{v,q}^\delta f(z))'}{\Psi_{v,q}^\delta f(z)} - 1 \right) \prec \Phi(z), \quad z \in \mathbb{D}.$$

By specializing the parameters, suitably we deduce the following new subclasses based on Bessel functions , which are not yet been studied.

Definition 3. For $\mu = 0, \lambda \geq 0$ and let $\Phi(z)$ be given in Definition 1 and $f \in \mathcal{A}$ is in the class $\mathcal{G}_{0\lambda}^{v, \delta, q}(\Phi) = \mathcal{M}_\lambda^{v, \delta, q}(\Phi)$ if

$$\lambda \left(1 + \frac{z(\Psi_{v,q}^\delta f(z))''}{(\Psi_{v,q}^\delta f(z))'} \right) + (1 - \lambda) \frac{z(\Psi_{v,q}^\delta f(z))'}{\Psi_{v,q}^\delta f(z)} \prec \Phi(z), \quad z \in \mathbb{D}.$$

Note that

$$\mathcal{M}_0^{v, \delta, q}(\Phi) \equiv \mathcal{S}^{v, \delta, q}(\Phi) \quad \text{and} \quad \mathcal{M}_1^{v, \delta, q}(\Phi) \equiv \mathcal{C}^{v, \delta, q}(\Phi).$$

Definition 4. For $\mu \geq 0, \lambda = 0$ and let $\Phi(z)$ be in Definition 1 and $f \in \mathcal{A}$ is in the class $\mathcal{G}_{\mu, 0}^{v, \delta, q}(\Phi) = \mathcal{B}_\mu^{v, \delta, q}(\Phi)$ if

$$(\Psi_{v,q}^\delta f(z))' \left(\frac{\Psi_{v,q}^\delta f(z)}{z} \right)^{\mu-1} \prec \Phi(z), \quad z \in \mathbb{D}.$$

Definition 5. For $\mu = 1, \lambda = 0$ and let $\Phi(z)$ be in Definition 1 and $f \in \mathcal{A}$ is in the class $\mathcal{G}_{1, 0}^{v, \delta, q}(\Phi) = \mathcal{R}^{v, \delta, q}(\Phi)$ if

$$(\Psi_{v,q}^\delta f(z))' \prec \Phi(z), \quad z \in \mathbb{D}.$$

To prove our main result, we need the following lemmas:

Lemma 1.[12] If $\varpi \in \mathcal{P}$ and given by

$$\varpi(z) = 1 + c_1z + c_2z^2 + \dots \quad (19)$$

then $|c_j| \leq 2$ for all $j \geq 1$, and the result is best possible for $\phi_1(z) = \frac{1 + \eta z}{1 - \eta z}, |\eta| = 1$.

Lemma 2.[9] If $\varpi(z) \in \mathcal{P}$ and given by (19) then

$$|c_2 - \vartheta c_1^2| \leq \begin{cases} -4\vartheta + 2, & \text{if } \vartheta \leq 0, \\ 2, & \text{if } 0 \leq \vartheta \leq 1, \\ 4\vartheta - 2, & \text{if } \vartheta \geq 1. \end{cases}$$

When $\vartheta < 0$ or $\vartheta > 1$, the equality holds if and only if $p_1(z) = \frac{1+z}{1-z}$ or one of its rotations. If $0 < \vartheta < 1$, then the above upper bound is sharp, for $p_2(z) = \frac{1+z^2}{1-z^2}$ or one of its rotations. If $\vartheta = 0$, the above upper bound is sharp, for

$$p_3(z) = \left(\frac{1}{2} + \frac{\eta}{2}\right) \frac{1+z}{1-z} + \left(\frac{1}{2} - \frac{\eta}{2}\right) \frac{1-z}{1+z} \quad (0 \leq \eta \leq 1)$$

or one of its rotations. If $\vartheta = 1$, the equality holds if and only if p_1 is the reciprocal of one of the functions such that the equality holds in the case of $\vartheta = 0$.

Although the above upper bound is sharp, when $0 < \vartheta < 1$, it can be improved as follows:

$$|c_2 - \vartheta c_1^2| + \vartheta |c_1|^2 \leq 2 \quad (0 < \vartheta \leq 1/2)$$

and

$$|c_2 - \vartheta c_1^2| + (1 - \vartheta) |c_1|^2 \leq 2 \quad (1/2 < \vartheta \leq 1).$$

We also need the following:

Lemma 3. [13] If $\omega(z) = 1 + c_1z + c_2z^2 + \dots$ is a member of \mathcal{P} , then

$$|c_2 - \vartheta c_1^2| \leq 2 \max(1, |2\vartheta - 1|).$$

The result is sharp for the functions

$$p(z) = \frac{1+z^2}{1-z^2}, \quad p(z) = \frac{1+z}{1-z} \quad (z \in \mathbb{D}).$$

2 Coefficient Estimates and Fekete-Szegő inequality

To start with in this section we determine the initial Coefficient estimates a_2 and a_3 . Unless otherwise stated, we let the following in our study:

$$\Theta_2 = \frac{(-1/4)\Gamma(v+1)}{\Gamma(2+v)} \frac{\Gamma_q(2+\delta)}{[1]_q! \Gamma_q(1+\delta)}$$

and

$$\Theta_3 = \frac{(-1/4)^2 \Gamma(v+1)}{2\Gamma(3+v)} \frac{\Gamma_q(3+\delta)}{[2]_q! \Gamma_q(1+\delta)}. \tag{20}$$

Theorem 1. Let $\mu \geq 0$ and $\lambda \geq 0$ and μ a real number and $\Phi(z)$ be given by (18). If $f(z)$ given by (1) belongs to $\mathcal{G}_{\mu,\lambda}^{v,\delta,q}(\Phi)$, then

$$|a_2| \leq \left| \frac{B_1}{\tau \Theta_2} \right|,$$

$$|a_3| \leq \left| \frac{B_1}{2\xi \Theta_3} \right| - \frac{B_2}{B_1} + \frac{B_1 \Lambda}{\tau^2} \Big|,$$

where $\Lambda = \frac{1}{2}(\mu^2 + \mu) - (\mu + 3)\lambda - 1$;

$$\tau := (1 + \mu)(1 + \lambda), \quad \text{and} \quad \xi := (\mu + 2)(1 + 2\lambda). \tag{21}$$

These results are sharp.

Proof. If $f \in \mathcal{G}_{\mu,\lambda}^{v,\delta,q}(\Phi)$, then there is a Schwarz function $\omega(z)$, analytic in \mathbb{D} with $\omega(0) = 0$ and $|\omega(z)| < 1$ in \mathbb{D} such that

$$\frac{z(\Psi_{v,q}^\delta f(z))'}{\Psi_{v,q}^\delta f(z)} \left(\frac{\Psi_{v,q}^\delta f(z)}{z} \right)^\mu + \lambda \left[1 + \frac{z(\Psi_{v,q}^\delta f(z))''}{(\Psi_{v,q}^\delta f(z))'} - \frac{z(\Psi_{v,q}^\delta f(z))'}{\Psi_{v,q}^\delta f(z)} \right] + \mu \left(\frac{z(\Psi_{v,q}^\delta f(z))'}{\Psi_{v,q}^\delta f(z)} - 1 \right) \prec \Phi(z) = \Phi(\omega(z)). \tag{22}$$

Define the function $p_1(z)$ by

$$p_1(z) := \frac{1 + \omega(z)}{1 - \omega(z)} = 1 + c_1z + c_2z^2 + \dots \tag{23}$$

Since $\omega(z)$ is a Schwarz function, we see that $\Re(p_1(z)) > 0$ and $p_1(0) = 1$. Let us define the function $p(z)$ by

$$p(z) := \frac{z(\Psi_{v,q}^\delta f(z))'}{\Psi_{v,q}^\delta f(z)} \left(\frac{\Psi_{v,q}^\delta f(z)}{z} \right)^\mu + \lambda \left[1 + \frac{z(\Psi_{v,q}^\delta f(z))''}{(\Psi_{v,q}^\delta f(z))'} - \frac{z(\Psi_{v,q}^\delta f(z))'}{\Psi_{v,q}^\delta f(z)} \right] + \mu \left(\frac{z(\Psi_{v,q}^\delta f(z))'}{\Psi_{v,q}^\delta f(z)} - 1 \right) = 1 + b_1z + b_2z^2 + \dots \tag{24}$$

In view of (22), (23), (24), we have

$$p(z) = \Phi \left(\frac{p_1(z) - 1}{p_1(z) + 1} \right). \tag{25}$$

Using (23) in (25), we get,

$$b_1 = \frac{1}{2} B_1 c_1 \quad \text{and} \quad b_2 = \frac{1}{2} B_1 \left(c_2 - \frac{1}{2} c_1^2 \right) + \frac{1}{4} B_2 c_1^2.$$

A computation shows that

$$\frac{z(\Psi_{v,q}^\delta f(z))'}{\Psi_{v,q}^\delta f(z)} = 1 + \Theta_2 a_2 z + (2\Theta_3 a_3 - \Theta_2^2 a_2^2) z^2 + (3\Theta_4 a_4 + \Theta_2^3 a_2^3 - 3\Theta_3 \Theta_2 a_3 a_2) z^3 + \dots$$

Similarly we have

$$1 + \frac{z(\Psi_{v,q}^\delta f(z))''}{(\Psi_{v,q}^\delta f(z))'} = 1 + 2\Theta_2 a_2 z + (6\Theta_3 a_3 - 4\Theta_2^2 a_2^2) z^2 + \dots$$

An easy computation shows that

$$\begin{aligned} & \frac{z(\Psi_{v,q}^\delta f(z))'}{\Psi_{v,q}^\delta f(z)} \left(\frac{\Psi_{v,q}^\delta f(z)}{z} \right)^\mu \\ & + \lambda \left[1 + \frac{z(\Psi_{v,q}^\delta f(z))''}{(\Psi_{v,q}^\delta f(z))'} - \frac{z(\Psi_{v,q}^\delta f(z))'}{\Psi_{v,q}^\delta f(z)} \right. \\ & \left. + \mu \left(\frac{z(\Psi_{v,q}^\delta f(z))'}{\Psi_{v,q}^\delta f(z)} - 1 \right) \right] \\ & = 1 + (1 + \mu)(1 + \lambda)\Theta_2 a_2 z + (\mu + 2)(1 + 2\lambda)\Theta_3 a_3 z^2 \\ & \quad + \left(\frac{\mu^2 + \mu}{2} - (\mu + 3)\lambda - 1 \right) \Theta_2^2 a_2^2 z^2 + \dots \end{aligned}$$

In prospect of (24), we see that

$$\begin{aligned} b_1 &= (1 + \mu)(1 + \lambda)\Theta_2 a_2 \\ b_2 &= (\mu + 2)(1 + 2\lambda)\Theta_3 a_3 \\ & \quad + \left(\frac{\mu^2 + \mu}{2} - (\mu + 3)\lambda - 1 \right) \Theta_2^2 a_2^2 \\ &= (\mu + 2)(1 + 2\lambda)\Theta_3 a_3 + \Lambda \Theta_2^2 a_2^2 \end{aligned}$$

where $\Lambda = \frac{1}{2}(\mu^2 + \mu) - (\mu + 3)\lambda - 1$. Equivalently, we have

$$\begin{aligned} a_2 &= \frac{B_1 c_1}{2(1 + \mu)(1 + \lambda)\Theta_2} = \frac{B_1 c_1}{2\tau\Theta_2}, \tag{26} \\ a_3 &= \frac{B_1}{2(\mu + 2)(1 + 2\lambda)\Theta_3} \\ & \quad \times \left[c_2 - \frac{1}{2} \left(1 - \frac{B_2}{B_1} + \frac{B_1\Lambda}{[(1 + \mu)(1 + \lambda)]^2} \right) c_1^2 \right] \\ &= \frac{B_1}{2\xi\Theta_3} \left[c_2 - \frac{1}{2} \left(1 - \frac{B_2}{B_1} + \frac{B_1\Lambda}{\tau^2} \right) c_1^2 \right]. \tag{27} \end{aligned}$$

From(26) and applying Lemma 1., we get

$$|a_2| \leq \left| \frac{B_1}{\tau\Theta_2} \right|.$$

From (27), by using the estimate

$$|c_2 - \vartheta c_1^2| \leq 2 \max\{1, |2\vartheta - 1|\}$$

where $\vartheta = \frac{1}{2} \left(1 - \frac{B_2}{B_1} + \frac{B_1\Lambda}{\tau^2} \right)$, given in Lemma 3 we have

$$\begin{aligned} |a_3| &\leq \frac{B_1}{(\mu + 2)(1 + 2\lambda)\Theta_3} \\ & \quad \times \max\{1, |2 \times \frac{1}{2} \left(1 - \frac{B_2}{B_1} + \frac{B_1\Lambda}{\tau^2} \right) - 1|\} \\ &= \frac{B_1}{\xi\Theta_3} \max\{1, \left| -\frac{B_2}{B_1} + \frac{B_1\Lambda}{\tau^2} \right|\} \end{aligned}$$

where τ, ξ are as assumed in (21) and $\Lambda = \frac{1}{2}(\mu^2 + \mu) - (\mu + 3)\lambda - 1$.

Using Lemma 2, we prove the following:

Theorem 2. Let $\mu \geq 0$ and $\lambda \geq 0$ and v a real number and $\Phi(z)$ be given by (18). If $f \in \mathcal{G}_{\mu,\lambda}^{v,\delta,q}(\Phi)(z)$, then

$$|a_3 - va_2^2| \leq \begin{cases} \frac{B_1}{\xi\Theta_3} \left(\frac{B_2}{B_1} - \frac{B_1\Lambda}{\tau^2} - \frac{v\xi B_1\Theta_3}{\tau^2\Theta_2^2} \right), & \text{if } v \leq \sigma_1, \\ \frac{B_1}{\Theta_3\xi}, & \text{if } \sigma_1 \leq v \leq \sigma_2, \\ \frac{B_1}{\xi\Theta_3} \left(-\frac{B_2}{B_1} + \frac{B_1\Lambda}{\tau^2} + \frac{v\xi B_1\Theta_3}{\tau^2\Theta_2^2} \right), & \text{if } v \geq \sigma_2, \end{cases}$$

where, for convenience,

$$\begin{aligned} \sigma_1 &:= \frac{\tau^2\Theta_2^2}{\xi B_1\Theta_3} \left(-1 + \frac{B_2}{B_1} - \frac{B_1\Lambda}{\tau^2} \right), \\ \sigma_2 &:= \frac{\tau^2\Theta_2^2}{\xi B_1\Theta_3} \left(1 + \frac{B_2}{B_1} - \frac{B_1\Lambda}{\tau^2} \right), \\ \sigma_3 &:= \frac{\tau^2\Theta_2^2}{\xi B_1\Theta_3} \left(\frac{B_2}{B_1} - \frac{B_1\Lambda}{\tau^2} \right), \end{aligned}$$

also τ, ξ are as defined in (21) and $\Lambda = \frac{1}{2}(\mu^2 + \mu) - (\mu + 3)\lambda - 1$.

Further, if $\sigma_1 \leq v \leq \sigma_3$, then

$$\begin{aligned} |a_3 - va_2^2| + \frac{\tau^2\Theta_2^2}{\xi B_1\Theta_3} \left(1 - \frac{B_2}{B_1} + \frac{B_1\Lambda}{\tau^2} + \frac{v\xi B_1\Theta_3}{\tau^2\Theta_2^2} \right) |a_2|^2 \\ \leq \frac{B_1}{\xi\Theta_3}. \end{aligned}$$

If $\sigma_3 \leq v \leq \sigma_2$, then

$$\begin{aligned} |a_3 - va_2^2| + \frac{\tau^2\Theta_2^2}{\xi B_1\Theta_3} \left(1 + \frac{B_2}{B_1} - \frac{B_1\Lambda}{\tau^2} - \frac{v\xi B_1\Theta_3}{\tau^2\Theta_2^2} \right) |a_2|^2 \\ \leq \frac{B_1}{\xi\Theta_3}. \end{aligned}$$

These results are sharp.

Proof. From (26) and (27), we have

$$\begin{aligned} a_3 - va_2^2 &= \frac{B_1}{2\xi\Theta_3} \left[c_2 - \frac{c_1^2}{2} \left(1 - \frac{B_2}{B_1} + \frac{B_1\Lambda}{\tau^2} \right) \right] \\ & \quad - \frac{c_1^2}{4} \frac{vB_1^2}{(\tau\Theta_2)^2} \\ &= \frac{B_1}{2\xi\Theta_3} \left[c_2 - \frac{c_1^2}{2} \left(1 - \frac{B_2}{B_1} + \frac{B_1\Lambda}{\tau^2} + \frac{v\xi B_1\Theta_3}{(\tau\Theta_2)^2} \right) \right]. \end{aligned}$$

Therefore, we have

$$a_3 - va_2^2 = \frac{B_1}{2\xi\Theta_3} (c_2 - vc_1^2)$$

where

$$v := \frac{1}{2} \left(1 - \frac{B_2}{B_1} + \frac{B_1\Lambda}{\tau^2} + \frac{v\xi B_1}{(\tau\Theta_2)^2} \right).$$

By applying Lemma 2 we get the desired result. To prove the bounds are sharp, we define the functions K_{Φ_n} ($n = 2, 3, \dots$) with $K_{\Phi_n}(0) = 0 = [K_{\Phi_n}]'(0) - 1$, by

$$\begin{aligned} & \frac{z(\Psi_{v,q}^\delta K_{\Phi_n}(z))'}{\Psi_{v,q}^\delta K_{\Phi_n}(z)} \left(\frac{\Psi_{v,q}^\delta K_{\Phi_n}(z)}{z} \right)^\mu \\ & + \lambda \left[1 + \frac{z(\Psi_{v,q}^\delta K_{\Phi_n}(z))''}{(\Psi_{v,q}^\delta K_{\Phi_n}(z))'} - \frac{z(\Psi_{v,q}^\delta K_{\Phi_n}(z))'}{\Psi_{v,q}^\delta K_{\Phi_n}(z)} \right. \\ & \left. + \mu \left(\frac{z(\Psi_{v,q}^\delta K_{\Phi_n}(z))'}{\Psi_{v,q}^\delta K_{\Phi_n}(z)} - 1 \right) \right] = \Phi(z^{n-1}) \end{aligned}$$

and the functions F_ρ and G_ρ ($0 \leq \rho \leq 1$), respectively, with $F_\rho(0) = 0 = F_\rho'(0) - 1$ and $G_\rho(0) = 0 = G_\rho'(0) - 1$ by

$$\begin{aligned} & \frac{z(\Psi_{v,q}^\delta F_\rho(z))'}{\Psi_{v,q}^\delta F_\rho(z)} \left(\frac{\Psi_{v,q}^\delta F_\rho(z)}{z} \right)^\mu \\ & + \lambda \left[1 + \frac{z(\Psi_{v,q}^\delta F_\rho(z))''}{(\Psi_{v,q}^\delta F_\rho(z))'} - \frac{z(\Psi_{v,q}^\delta F_\rho(z))'}{\Psi_{v,q}^\delta F_\rho(z)} \right. \\ & \left. + \mu \left(\frac{z(\Psi_{v,q}^\delta F_\rho(z))'}{\Psi_{v,q}^\delta F_\rho(z)} - 1 \right) \right] = \Phi \left(\frac{z(z+\rho)}{1+\rho z} \right), \end{aligned}$$

and

$$\begin{aligned} & \frac{z(\Psi_{v,q}^\delta G_\rho(z))'}{\Psi_{v,q}^\delta G_\rho(z)} \left(\frac{\Psi_{v,q}^\delta G_\rho(z)}{z} \right)^\mu \\ & + \lambda \left[1 + \frac{z(\Psi_{v,q}^\delta G_\rho(z))''}{(\Psi_{v,q}^\delta G_\rho(z))'} - \frac{z(\Psi_{v,q}^\delta G_\rho(z))'}{\Psi_{v,q}^\delta G_\rho(z)} \right. \\ & \left. + \mu \left(\frac{z(\Psi_{v,q}^\delta G_\rho(z))'}{\Psi_{v,q}^\delta G_\rho(z)} - 1 \right) \right] = \Phi \left(-\frac{z(z+\rho)}{1+\rho z} \right), \end{aligned}$$

respectively.

Clearly the functions $K_{\Phi_n}, F_\rho, G_\rho \in \mathcal{G}_{\mu,\lambda}^{v,\delta,q}(\Phi)$. Also we write $K_\Phi := K_{\Phi_2}$.

If $\mu < \sigma_1$ or $\mu > \sigma_2$, then the bounds are sharp if and only if f is K_Φ or one of its rotations. When $\sigma_1 < \mu < \sigma_2$, then the bounds are sharp if and only if f is K_{Φ_3} or one of its rotations. If $\mu = \sigma_1$ then the bounds are sharp if and only if f is F_ρ or one of its rotations. If $\mu = \sigma_2$ then the bounds are sharp if and only if f is G_ρ or one of its rotations.

By making use of Lemma 3, we immediately obtain the following:

Theorem 3. Let $\mu \geq 0$ and $\lambda \geq 0$ further, let $\Phi(z)$ be of the form (18). If $f \in \mathcal{G}_{\mu,\lambda}^{v,\delta,q}(\Phi)$, then for complex v , we have

$$|a_3 - va_2^2| = \frac{B_1}{\xi \Theta_3} \max \left\{ 1, \left| -\frac{B_2}{B_1} + \frac{B_1 \Lambda}{\tau^2} + \frac{v \xi \Theta_3}{\tau^2 \Theta_2^2} B_1 \right| \right\}$$

where ξ, τ are as assumed in (21) and $\Lambda = \frac{1}{2}(\mu^2 + \mu) - (\mu + 3)\lambda - 1$. The result is sharp.

3 Coefficient inequalities for the function

$$f^{-1} \in \mathcal{G}_{\mu,\lambda}^{v,\delta,q}(\Phi)$$

Theorem 4. If $f \in \mathcal{G}_{\mu,\lambda}^{v,\delta,q}(\Phi)$ and $f^{-1}(w) = w + \sum_{n=2}^{\infty} d_n w^n$ is the inverse function of f with ($|w| < r_0(f); r_0(f) \geq \frac{1}{4}$) the Koebe domain of the class $f \in \mathcal{G}_{\mu,\lambda}^{v,\delta,q}(\Phi)$, then for any complex number v , we have

$$\begin{aligned} |d_3 - vd_2^2| & \leq \frac{B_1}{\xi \Theta_3} \\ & \times \max \left\{ 1, \left| -\frac{B_2}{B_1} + \frac{B_1 \Lambda}{\tau^2} + \frac{(2-v)B_1 \xi \Theta_3}{\tau^2 \Theta_2^2} \right| \right\}. \end{aligned} \tag{28}$$

where τ, ξ are as assumed in (21) and $\Lambda = \frac{1}{2}(\mu^2 + \mu) - (\mu + 3)\lambda - 1$.

Proof. As

$$f^{-1}(w) = w + \sum_{n=2}^{\infty} d_n w^n \tag{29}$$

is the inverse function of f , it can be seen that

$$f^{-1}(f(z)) = f\{f^{-1}(z)\} = z. \tag{30}$$

From equations (1) and (30), we get

$$f^{-1}(z + \sum_{n=2}^{\infty} a_n z^n) = z. \tag{31}$$

From (30) and (31), one can obtain

$$z + (a_2 + d_2)z^2 + (a_3 + 2a_2d_2 + d_3)z^3 + \dots = z. \tag{32}$$

By equating corresponding coefficients, of (32), we have

$$d_2 = -a_2 \tag{33}$$

$$d_3 = 2a_2^2 - a_3. \tag{34}$$

From relations (26),(27),(33) and (34)

$$d_2 = -\frac{B_1 c_1}{2(1+\mu)(1+\lambda)\Theta_2} = -\frac{B_1 c_1}{2\tau\Theta_2}; \tag{35}$$

$$\begin{aligned} d_3 & = \frac{B_1^2 c_1^2}{2\tau^2 \Theta_2^2} - \frac{B_1}{2\xi \Theta_3} \left(c_2 - \frac{1}{2} \left(1 - \frac{B_2}{B_1} + \frac{B_1 \Lambda}{\tau^2} \right) c_1^2 \right); \\ & = \frac{B_1}{2\xi \Theta_3} \left[-c_2 + \frac{c_1^2}{2} \left(1 - \frac{B_2}{B_1} + \frac{B_1 \Lambda}{\tau^2} + \frac{2B_1 \xi \Theta_3}{\tau^2 \Theta_2^2} \right) \right]; \\ & = \frac{-B_1}{2\xi \Theta_3} \left[c_2 - \frac{c_1^2}{2} \left(1 - \frac{B_2}{B_1} + \frac{B_1 \Lambda}{\tau^2} + \frac{2B_1 \xi \Theta_3}{\tau^2 \Theta_2^2} \right) \right]; \end{aligned} \tag{36}$$

where τ, ξ are as assumed in (21) and $\Lambda = \frac{1}{2}(\mu^2 + \mu) - (\mu + 3)\lambda - 1$. For any complex number v , consider

$$\begin{aligned} d_3 - vd_2^2 & = -\frac{B_1}{2\xi \Theta_3} \\ & \times \left[c_2 - \frac{c_1^2}{2} \left(1 - \frac{B_2}{B_1} + \frac{B_1 \Lambda}{\tau^2} + \frac{(2-v)B_1 \xi \Theta_3}{\tau^2 \Theta_2^2} \right) \right]. \end{aligned} \tag{37}$$

Taking modulus on both sides and by applying Lemma 3 on the right hand side of (37), one can obtain the result as in (28). Hence this completes the proof.

4 Application to Functions Defined by Poisson distribution

A variable χ is said to be Poisson distributed if it takes the values $0, 1, 2, 3, \dots$ with probabilities $e^{-\kappa}, m \frac{e^{-\kappa}}{1!}, \kappa^2 \frac{e^{-\kappa}}{2!}, \kappa^3 \frac{e^{-\kappa}}{3!}, \dots$ respectively, where κ is called the parameter. Thus

$$P(\chi = r) = \frac{\kappa^r e^{-\kappa}}{r!}, \quad r = 0, 1, 2, 3, \dots$$

In [14], Porwal introduced a power series whose coefficients are probabilities of Poisson distribution

$$\mathcal{J}(\kappa, z) = z + \sum_{n=2}^{\infty} \frac{\kappa^{n-1}}{(n-1)!} e^{-\kappa} z^n, \quad z \in \mathbb{D},$$

where $\kappa > 0$. We note by the familiar ratio test that the radius of convergence of the above series is infinity. Lately, using the Hadamard product, Porwal [14] (see also, [15, 16, 17]) introduced a new linear operator $\mathcal{J}^\kappa(z) : \mathcal{A} \rightarrow \mathcal{A}$ defined by

$$\begin{aligned} \mathcal{J}^\kappa f &= \mathcal{J}(\kappa, z) * f(z) = z + \sum_{n=2}^{\infty} \frac{\kappa^{n-1}}{(n-1)!} e^{-\kappa} a_n z^n, \\ &= z + \sum_{n=2}^{\infty} \psi_n(\kappa) a_n z^n, \quad z \in \mathbb{D}, \end{aligned}$$

where

$$\psi_n = \psi_n(\kappa) = \frac{\kappa^{n-1}}{(n-1)!} e^{-\kappa}.$$

In particular

$$\psi_2 = \kappa e^{-\kappa} \text{ and } \psi_3 = \frac{\kappa^2}{2} e^{-\kappa} \tag{38}$$

We describe the class $\mathcal{G}_{\kappa, \mu, \lambda}^{v, \delta, q}(\Phi)$ as below:

$$\mathcal{G}_{\kappa, \mu, \lambda}^{v, \delta, q}(\Phi) := \{f \in \mathcal{A} \text{ and } \mathcal{J}^\kappa f \in \mathcal{G}_{\mu, \lambda}^{v, \delta, q}(\Phi)\}$$

where $\mathcal{G}_{\mu, \lambda}^{v, \delta, q}(\Phi)$ is given by Definition 2. We find the coefficient estimate for $f \in \mathcal{G}_{\kappa, \mu, \lambda}^{v, \delta, q}(\Phi)$, from the corresponding estimate for functions in the class $\mathcal{G}_{\mu, \lambda}^{v, \delta, q}(\Phi)$. Applying Theorem 2 and 3, for $\mathcal{J}^\kappa f = \mathcal{J}(\kappa, z) * f(z) = z + \psi_2 a_2 z^2 + \psi_3 a_3 z^3 + \dots$, we get the following Theorems 5 and 6 after a noticeable modification of the parameter v .

Theorem 5. Let $\mu \geq 0; \lambda \geq 0$ and $\Phi(z)$ be given by (18). If $f \in \mathcal{G}_{\kappa, \mu, \lambda}^{v, \delta, q}(\Phi)$, then for complex μ , we have

$$|a_3 - v a_2^2| = \frac{2B_1}{\xi \Theta_3 \kappa^2 e^{-\kappa}} \times \max \left\{ 1, \left| -\frac{B_2}{B_1} + \frac{B_1 \Lambda}{\tau^2} + \frac{v \xi B_1 \Theta_3}{2\tau^2 \Theta_2^2 e^{-\kappa}} \right| \right\}$$

where τ, ξ are as assumed in (21) and $\Lambda = \frac{1}{2}(\mu^2 + \mu) - (\mu + 3)\lambda - 1$. The result is sharp.

Theorem 6. Let $\mu \geq 0; \lambda \geq 0$ and v a real number. Further, let $\Phi(z)$ be given in (18). If $f \in \mathcal{G}_{\kappa, \mu, \lambda}^{v, \delta, q}(\Phi)$, then

$$|a_3 - v a_2^2| \leq \begin{cases} \frac{2B_1}{\xi \Theta_3 \kappa^2 e^{-\kappa}} \left(\frac{B_2}{B_1} - \frac{B_1 \Lambda}{\tau^2} - \frac{v \xi B_1 \Theta_3}{2\tau^2 \Theta_2^2 e^{-\kappa}} \right), & \text{if } v \leq \sigma_1, \\ \frac{2B_1}{\xi \Theta_3 \kappa^2 e^{-\kappa}}, & \text{if } \sigma_1 \leq v \leq \sigma_2, \\ \frac{2B_1}{\xi \Theta_3 \kappa^2 e^{-\kappa}} \left(-\frac{B_2}{B_1} + \frac{B_1 \Lambda}{\tau^2} + \frac{v \xi B_1 \Theta_3}{2\tau^2 \Theta_2^2 e^{-\kappa}} \right), & \text{if } v \geq \sigma_2, \end{cases}$$

where, for convenience,

$$\begin{aligned} \sigma_1 &= \frac{2\tau^2 \Theta_2^2 e^{-\kappa}}{\xi B_1 \Theta_3} \left(-1 + \frac{B_2}{B_1} - \frac{B_1 \Lambda}{\tau^2} \right), \\ \sigma_2 &= \frac{2\tau^2 \Theta_2^2 e^{-\kappa}}{\xi B_1 \Theta_3} \left(1 + \frac{B_2}{B_1} - \frac{B_1 \Lambda}{\tau^2} \right) \end{aligned}$$

$\Lambda = \frac{1}{2}(\mu^2 + \mu) - (\mu + 3)\lambda - 1$ and τ, ξ are as assumed in (21). These results are sharp.

Concluding Remarks

Suitably specializing the parameters μ and λ as stated in Definitions 3 to 5, in Theorems 2,3 and 4 one can easily state above result for the function classes defined in Definitions 3 to 5 related with Bessel Functions. Also, further fixing Φ as illustrated below:

1. For $(0 < \alpha \leq 1)$ and $-1 \leq B < A \leq 1$, taking the function Φ as

$$\begin{aligned} \Phi(z) &= \left(\frac{1 + Az}{1 + Bz} \right)^\alpha \\ &= 1 + \alpha(A - B)z \\ &\quad - \frac{\alpha}{2} [2B(A - B) + (1 - \alpha)(A - B)^2] z^2 + \dots \end{aligned}$$

which gives $B_1 = \alpha(A - B)$ and $B_2 = -\frac{\alpha}{2} [2B(A - B) + (1 - \alpha)(A - B)^2]$.

2. If we take $\alpha = 1$ and $-1 \leq B < A \leq 1$, then we have

$$\Phi(z) = \frac{1 + Az}{1 + Bz} = 1 + (A - B)z + B(A - B)z^2 + \dots \tag{39}$$

thus we have $B_1 = A - B$ and $B_2 = B(A - B)$.

3. By fixing $A = 1$ and $B = -1$ we have

$$\Phi(z) = \frac{1+z}{1-z} = 1 + 2z + 2z^2 + \dots \quad (40)$$

thus we have $B_1 = 2$ and $B_2 = 2$

4. Further for some $c \in (0, 1]$, taking

$$\Phi(z) = \sqrt{1+cz} = 1 + \frac{c}{2}z - \frac{c^2}{8}z^2 + \dots \quad (41)$$

then the class is said to be associated with the right-loop of the Cassinian Ovals [18]. In particular if $c = 1$ then the class is associated with right-half of the lemniscate of Bernoulli [19] is given by

$$\Phi(z) = \sqrt{1+z} = 1 + \frac{1}{2}z - \frac{1}{8}z^2 + \dots \quad (42)$$

which gives $B_1 = \frac{1}{2}$ and $B_2 = -\frac{1}{8}$.

5. Taking

$$\Phi(z) = z + \sqrt{1+z^2} = 1 + z + \frac{1}{2}z^2 - \frac{1}{8}z^4 + \dots \quad (43)$$

which gives $B_1 = 1$ and $B_2 = \frac{1}{2}$,

then the class is said to be associated with the right crescent [20].

6. Again by taking

$$\Phi(z) = 1 + \frac{4}{3}z + \frac{2}{3}z^2 \quad (44)$$

which gives $B_1 = \frac{4}{3}$ and $B_2 = \frac{2}{3}$,

then the class is said to be associated with the cardioid [21].

7. By taking $\Phi(z) = 1 + \frac{8}{\pi^2}z + \frac{16}{3\pi^2}z^2$ where f is a parabolic starlike function (see [22]) in conic regions, one can deduce the analogues results of above theorems, we left the proof as exercise to interested readers.

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Conflict of Interest The authors declare that they have no conflict of interest

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